

# Statistical analysis of beta decays and the effective value of $g_A$

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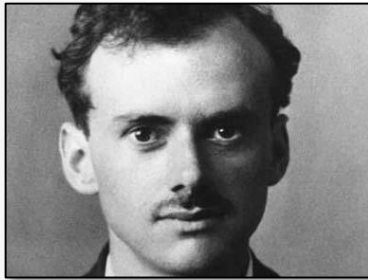
University College London

with J. Suhonen: Phys.Rev. C94 (2016) 5, 055501

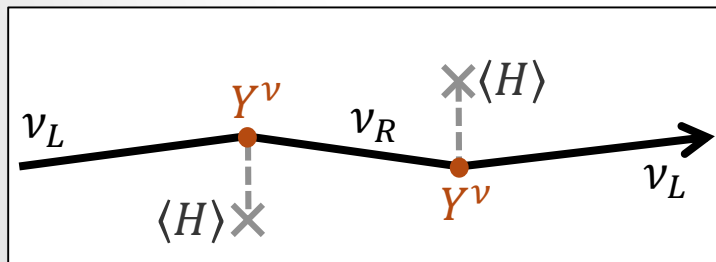
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# Dirac vs Majorana

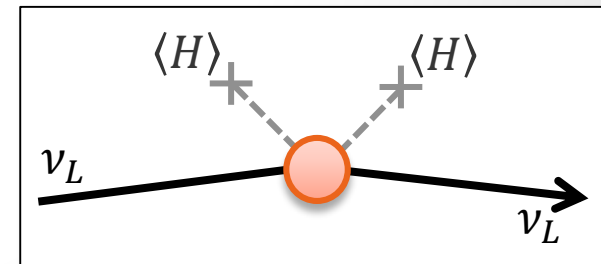
- ▶ Two possibilities to define fermion mass



Dirac mass analogous to other fermions but with  $m_\nu / \Lambda_{EW} \approx 10^{-12}$  couplings to Higgs



Majorana mass, using only a left-handed neutrino  
 → Lepton Number Violation



# Beta decays

▶ Single beta decay

$$(A, Z) \rightarrow (A, Z + 1) + e^- + \bar{\nu}_e$$

▶ Allowed double beta ( $2\nu\beta\beta$ ) decay

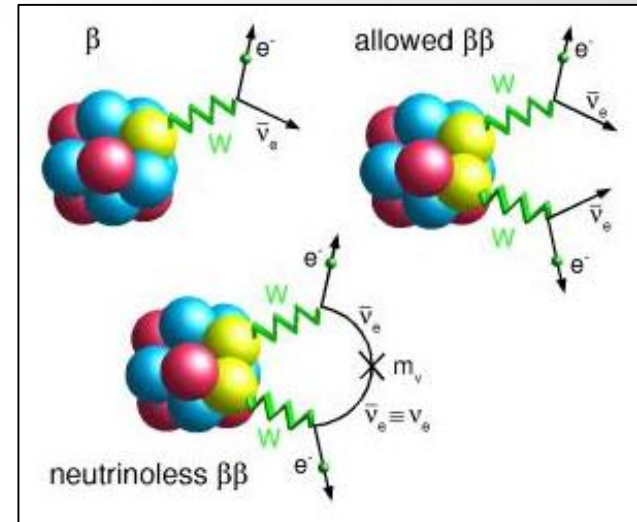
$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}_e$$

▶ Neutrinoless double beta ( $0\nu\beta\beta$ ) decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

- Violation of lepton number
- Mediated by Majorana neutrinos
- Variants

- $0\nu\beta^+\beta^+$ :  $(A, Z) \rightarrow (A, Z - 2) + 2e^+$
- $0\nu\beta^+EC$ :  $(A, Z) + e^- \rightarrow (A, Z - 2) + e^+$
- $0\nu ECEC$ :  $(A, Z) + 2e^- \rightarrow (A, Z - 2)$

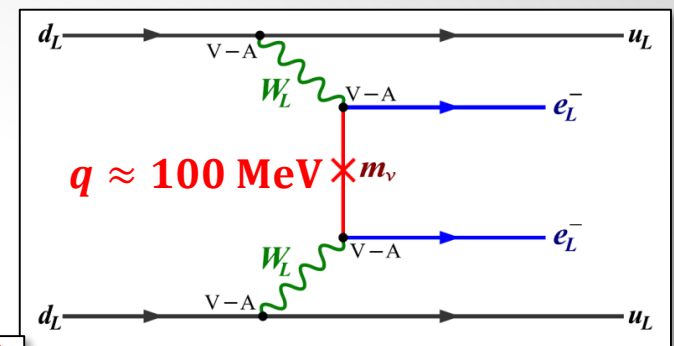


# $0\nu\beta\beta$

- ▶ Half-life  $T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$

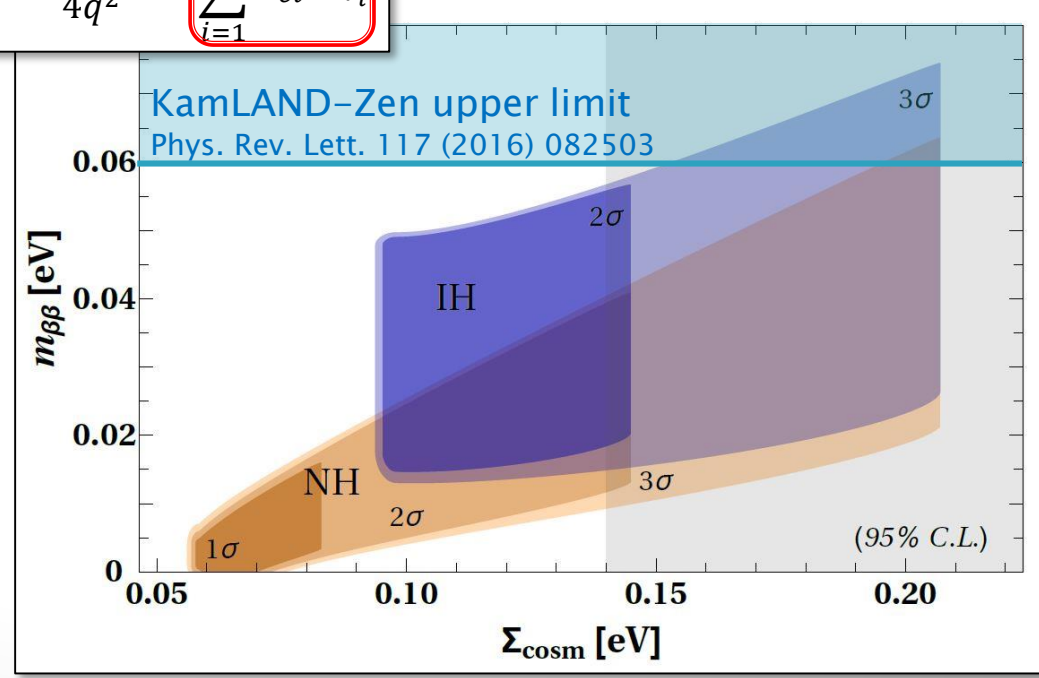
- ▶ Particle Physics

$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 U_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{q + \cancel{m}_{\nu_i}}{q^2 - m_{\nu_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4q^2} \sum_{i=1}^3 U_{ei}^2 m_{\nu_i}$$



- ▶ Atomic Physics
  - Leptonic phase space  $G^{0\nu}$
- ▶ Nuclear Physics
  - Nuclear transition matrix element  $M^{0\nu}$

$$\frac{10^{25} \text{yr}}{T_{1/2}} \approx \left( \frac{|m_{\beta\beta}|}{\text{eV}} \right)^2$$



Dell'Oro, Marcocci, Viel, Vissani, Adv.High Energy Phys. (2016) 2162659

# Nuclear Matrix Elements

- ▶ Hadronic current

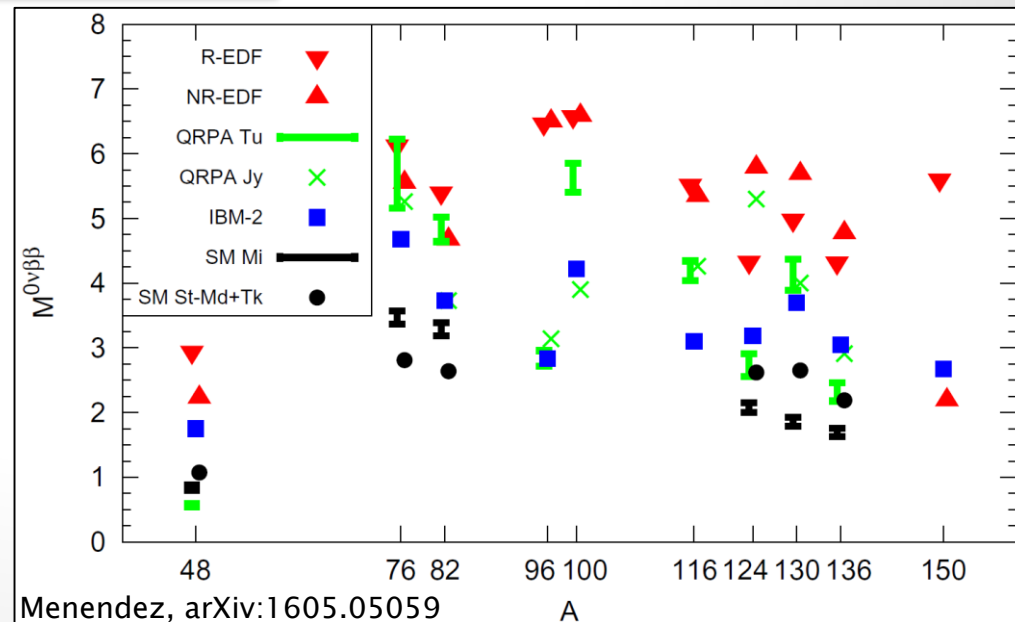
$$J^\mu(q) = g_V \gamma^\mu - g_A \gamma^\mu \gamma^5 + \frac{ig_M}{2m_N} \sigma^{\mu\nu} q_\nu - g_P \gamma^5 q^\mu$$

- ▶ Nuclear Matrix Element  $M^{0\nu}$

$$M^{0\nu} = g_A^2 \left( M_{GT} - \frac{g_V^2}{g_A^2} M_F + M_T \right)$$

IBM-3 MCSM  
IBM-1 NCSM  
BCS ISM IBM-1  
EDF NR-EDF IMSRG  
HFB QRPA CCEI  
GCM RPA  
R-EDF abinitio  
USDB IBM-4 IBM-2  
RQRPA

- Dependence on isotope and operator
- Many-body problem
- Factor 2 - 3 uncertainty between nuclear models



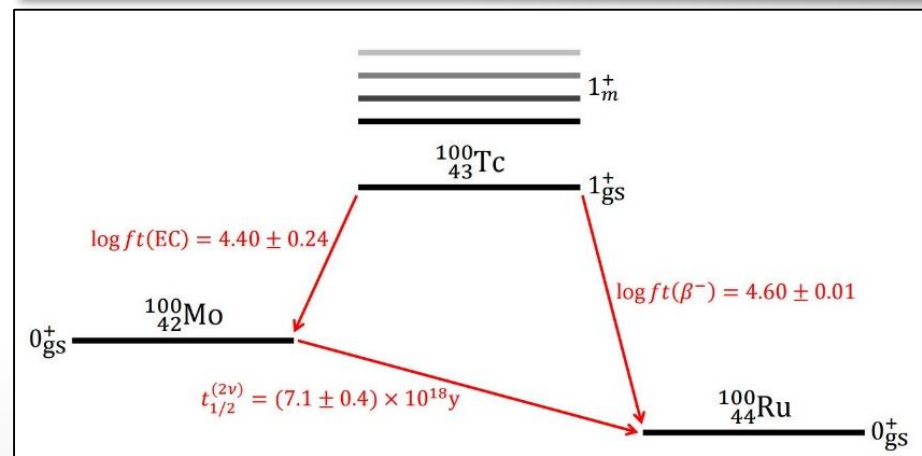
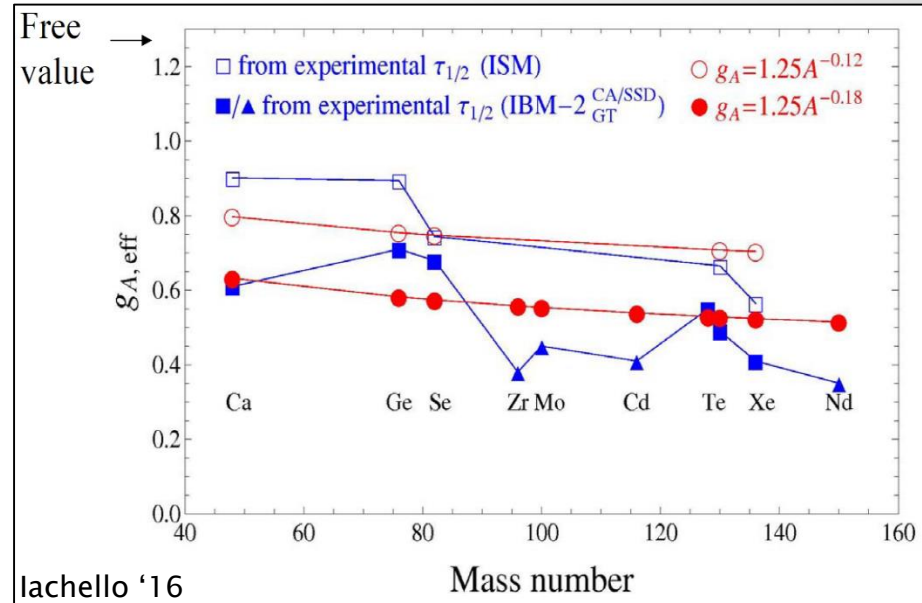
# Quenching of $g_A$ ?

▶ Nuclear matrix element

$$M^{0\nu} = g_A^2 \left( M_{GT} - \frac{g_V^2}{g_A^2} M_F + M_T \right)$$

▶ Axial-vector coupling  $g_A$

- Free nucleon:  $g_A \approx 1.27$
- Comparison of  $\beta$  and  $2\nu\beta\beta$  decay with theory:  
 $g_A \approx 0.6-0.8$
- If applicable to  $0\nu\beta\beta$ , strong reduction of sensitivity
- Genuine effect or short-coming of models?



## ▶ Calculation of single beta / EC and $2\nu\beta\beta$ matrix elements

J. Suhonen, O. Civitarese, Nucl. Phys. A 924 (2014) 1

- Woods–Saxon Potential with three different orbital spaces (A=62–80, 98 – 108, 110 – 142)
- Renormalized Bonn–A G–Matrix
- $g_{pp}$  and  $g_{ph}$  interaction parameters (per even–even system)
- $g_{ph}$  determined by location of Gamow–Teller Giant Resonance

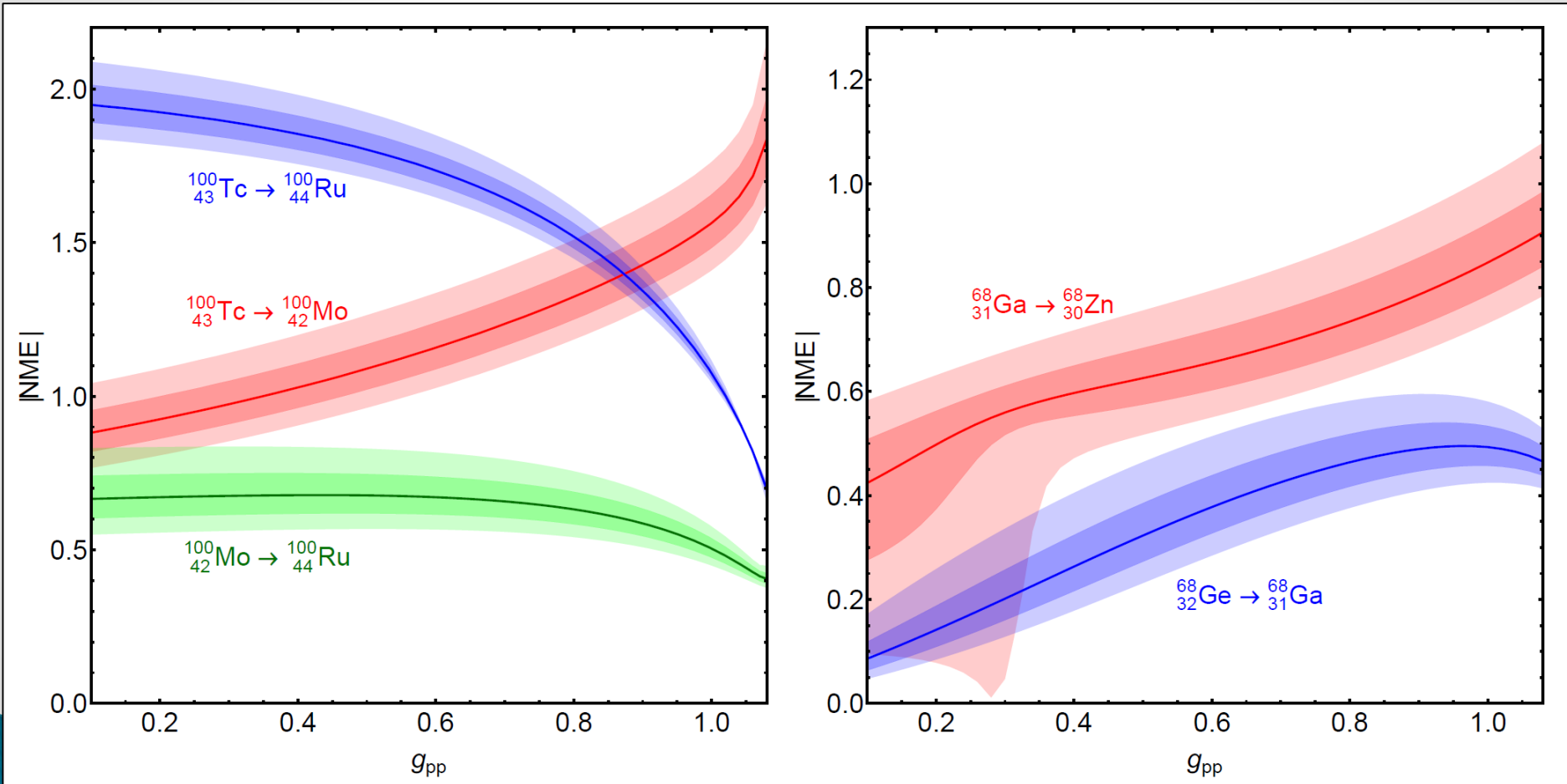
$$\Delta E_{GT} = E(1_{GTGR}^+) - E(0_{gs}^+) = \left[ 1.444 (Z + 1/2) A^{-1/3} - 30.0 (N - Z - 2) A^{-1} + 5.57 \right] \text{ MeV.}$$

only on average of theoretical / pheno centroids

→ include uncertainty of 15%

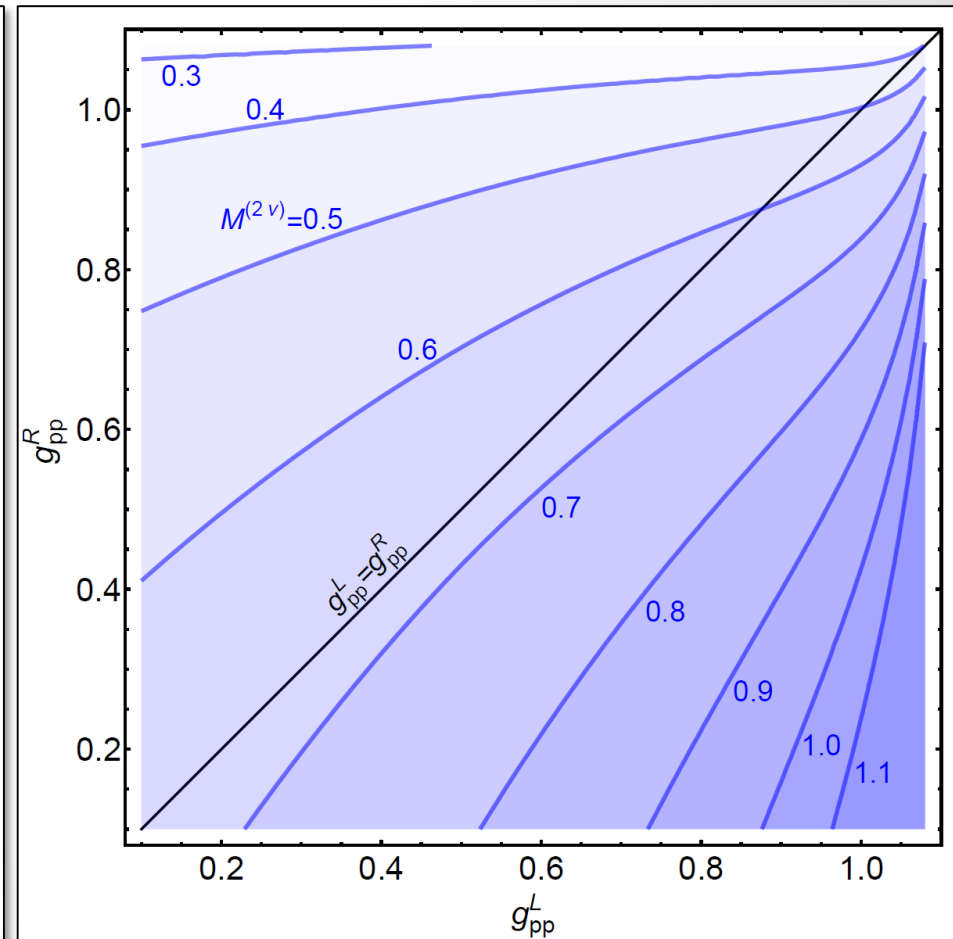
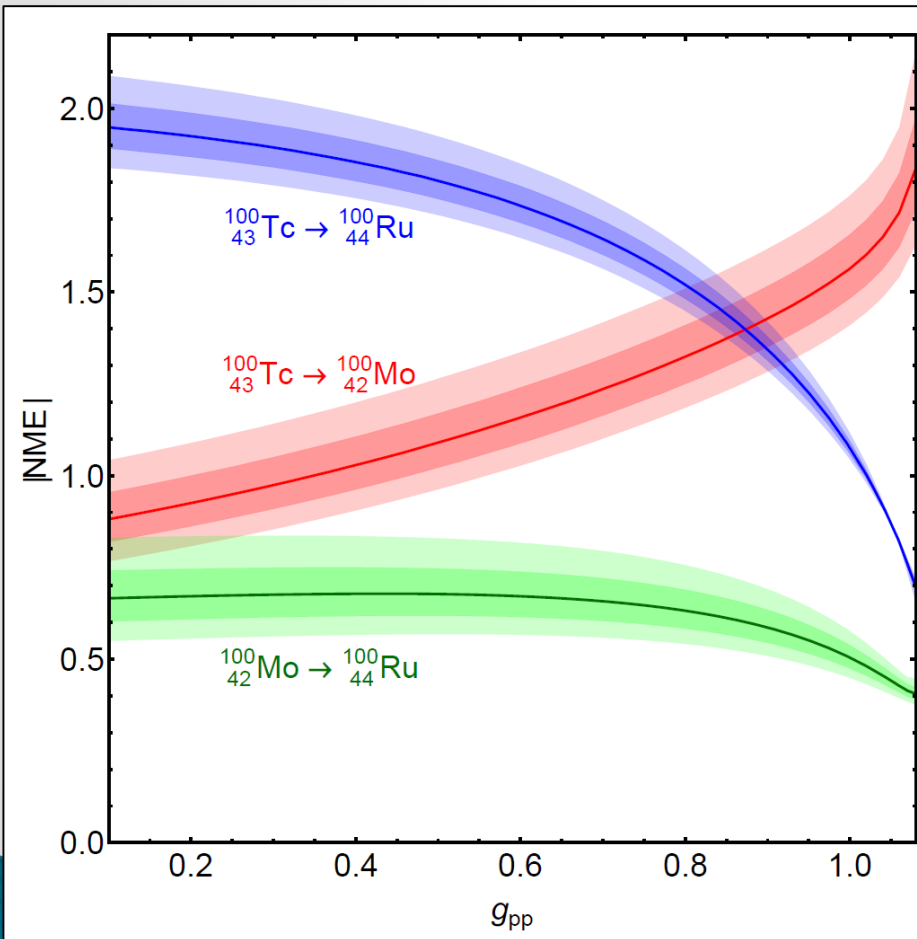
- $g_A$  as a free parameter per isobaric system

- ▶ Calculation of single beta / EC and  $2\nu\beta\beta$  matrix elements





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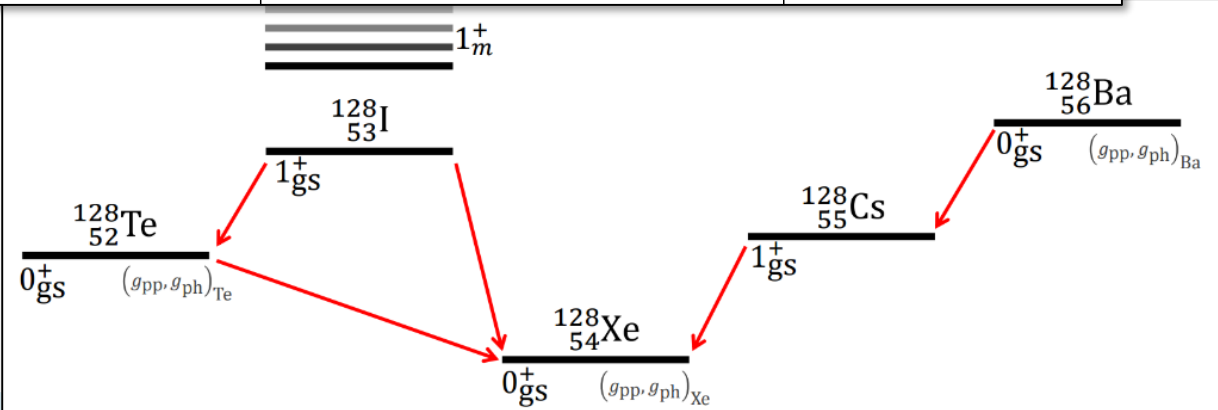
# Comparison with Beta Decay / EC Measurements

- ▶ Fit of model parameters over isobaric system
  - $g_A$
  - $g_{pp}^i$  (per even-even system)
  - $g_{ph}^i$  (per even-even system) as “nuisance” parameters
- ▶ Incorporate allowed  $gs \rightarrow gs$  GT beta decay / EC rates

$A$	$Z_0$	Triplet	$\log ft_L^{\text{exp}}$	$\log ft_R^{\text{exp}}$	$g_A^{\text{fit}}$	$g_{pp}^{\text{fit}}$
116	48	Cd(0 <sup>+</sup> ) ← In(1 <sup>+</sup> ) → Sn(0 <sup>+</sup> )	$4.4508 \pm 0.1160$	$4.6839 \pm 0.0025$	$0.84^{+0.08}_{-0.08}$	$0.65^{+0.07}_{-0.11}$
118	48	Cd(0 <sup>+</sup> ) → In(1 <sup>+</sup> ) → Sn(0 <sup>+</sup> )	$3.9218 \pm 0.0629$	$4.8147 \pm 0.0263$	$0.88^{+0.09}_{-0.07}$	$0.75^{+0.04}_{-0.09}$
118	49	In(1 <sup>+</sup> ) → Sn(0 <sup>+</sup> ) ← Sb(1 <sup>+</sup> )	$4.8147 \pm 0.0263$	$4.5152 \pm 0.0122$	$0.77^{+0.05}_{-0.06}$	$0.65^{+0.03}_{-0.04}$
118	50	Sn(0 <sup>+</sup> ) ← Sb(1 <sup>+</sup> ) ← Te(0 <sup>+</sup> )	$4.5152 \pm 0.0122$	$4.9749 \pm 0.0579$	$0.77^{+0.06}_{-0.05}$	$0.65^{+0.04}_{-0.14}$

$$\log_{10}(f_0 t_{1/2}[\text{s}]) = \log_{10}\left(\frac{6147}{B_{GT}}\right)$$

$$B_{GT} = \frac{g_A^2}{2J+1} |M_{GT}(g_{pp}, g_{ph})|^2$$



# Comparison with Beta Decay / EC Measurements

- ▶ Fit of model parameters over isobaric system
  - $g_A$
  - $g_{pp}^i$  (per even-even system)
  - $g_{ph}^i$  (per even-even system) as “nuisance” parameters
- ▶ Markov Chain Monte Carlo to find posterior parameter distribution  $p(g_A, g_{pp}^i, g_{ph}^i)$  based on fitness  $P = e^{-\chi^2/2}$ , e.g. triplet

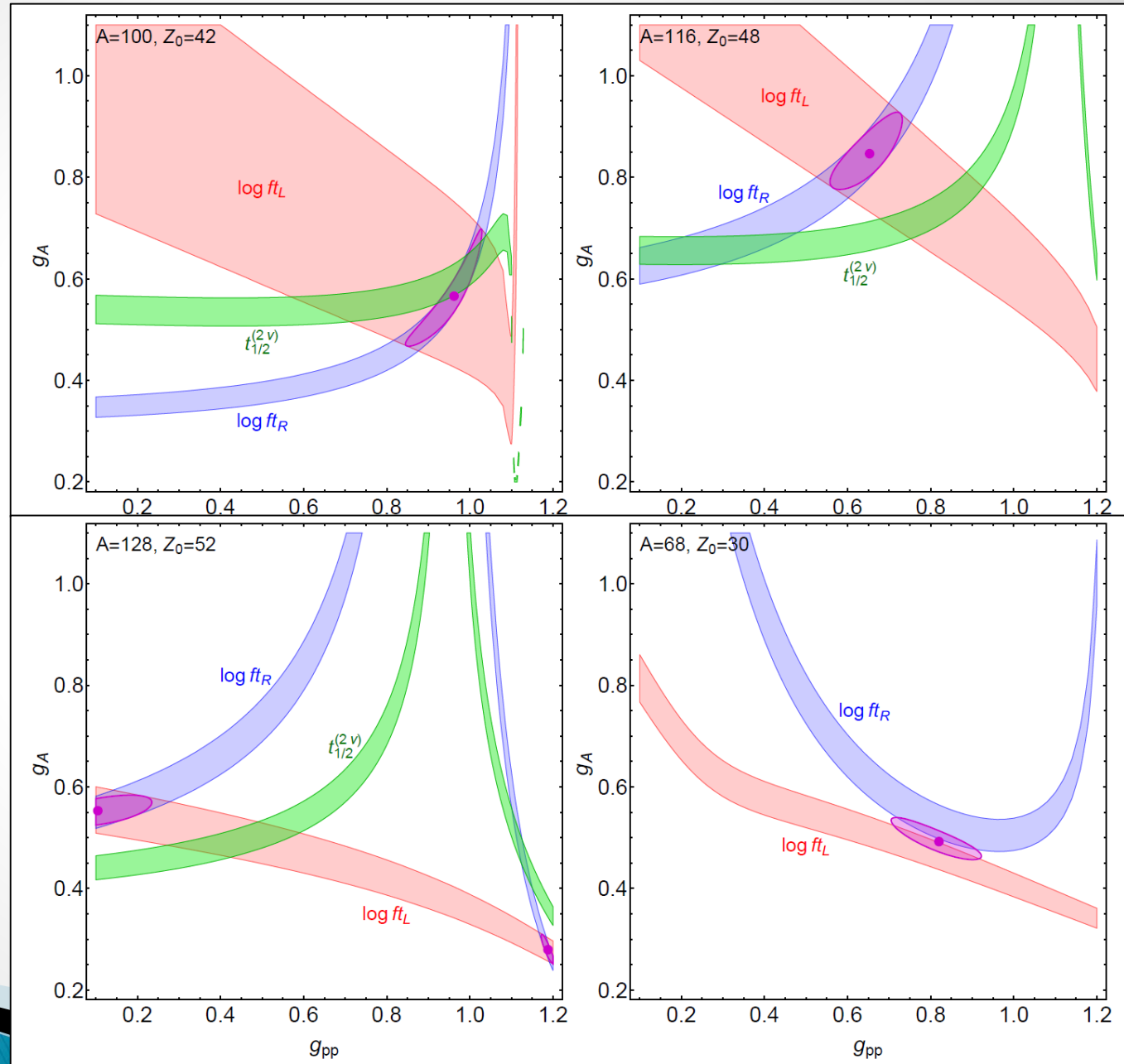
$$\chi^2 = \left( \frac{\log ft_L^{th}(g_A, g_{pp}^L, \gamma_{ph}^L) - \log ft_L^{exp}}{\delta \log ft_L^{exp}} \right)^2 + \left( \frac{\log ft_R^{th}(g_A, g_{pp}^R, \gamma_{ph}^R) - \log ft_R^{exp}}{\delta \log ft_R^{exp}} \right)^2 + \left( \frac{\gamma_{ph}^R - 1}{0.15} \right)^2 + \left( \frac{\gamma_{ph}^L - 1}{0.15} \right)^2$$

# Simplified Triplet Case

- ▶ Two parameters

$$g_A, g_{pp}$$

- ▶ See also  
A. Faessler et al.,  
J. Phys. G 35 (2008)  
075104



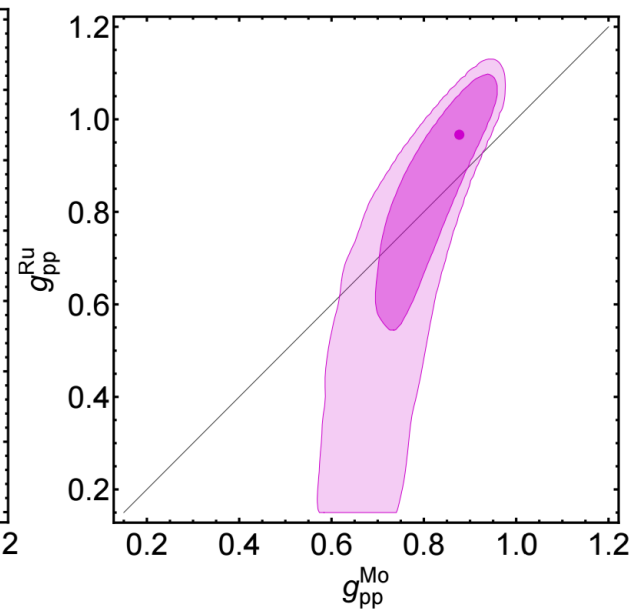
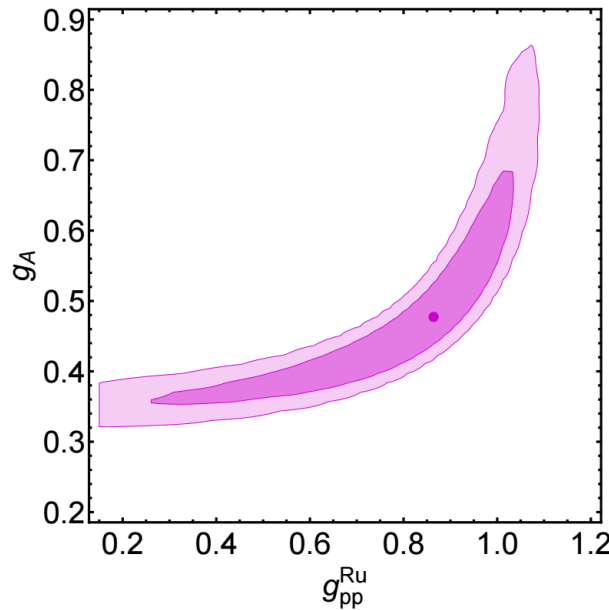
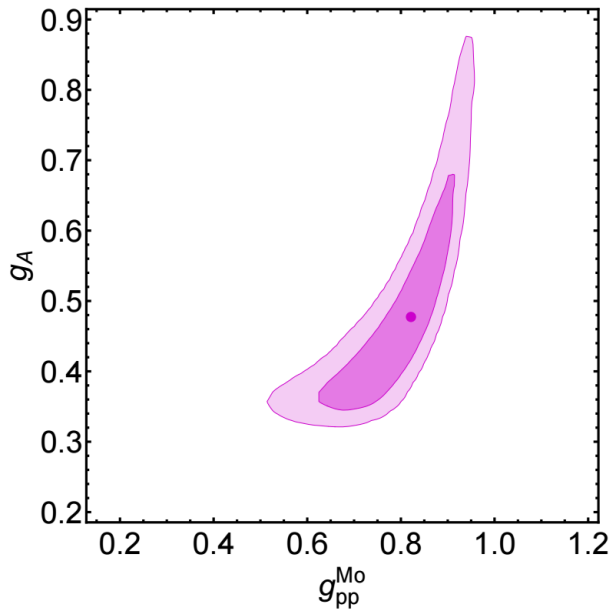
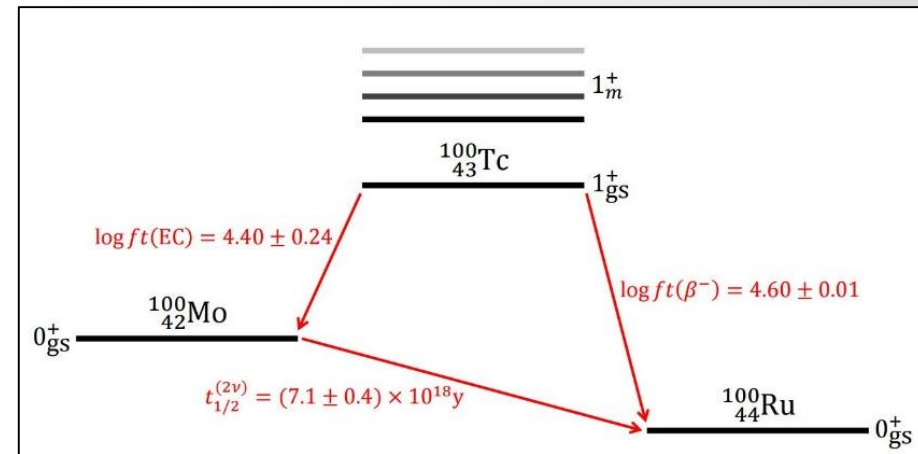
# Triplet Case

- ▶ Three parameters

$$g_A, g_{pp}^{\text{Mo}}, g_{pp}^{\text{Ru}}$$

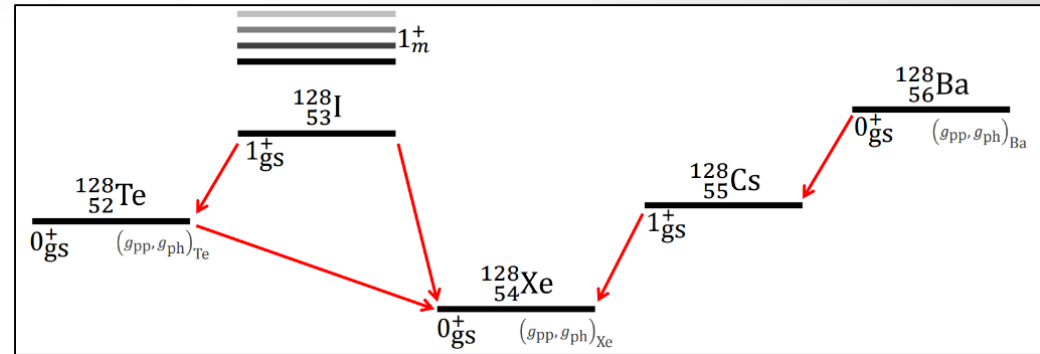
- ▶ Two nuisance

$$\text{parameters } \gamma_{ph}^{\text{Mo}}, \gamma_{ph}^{\text{Ru}}$$

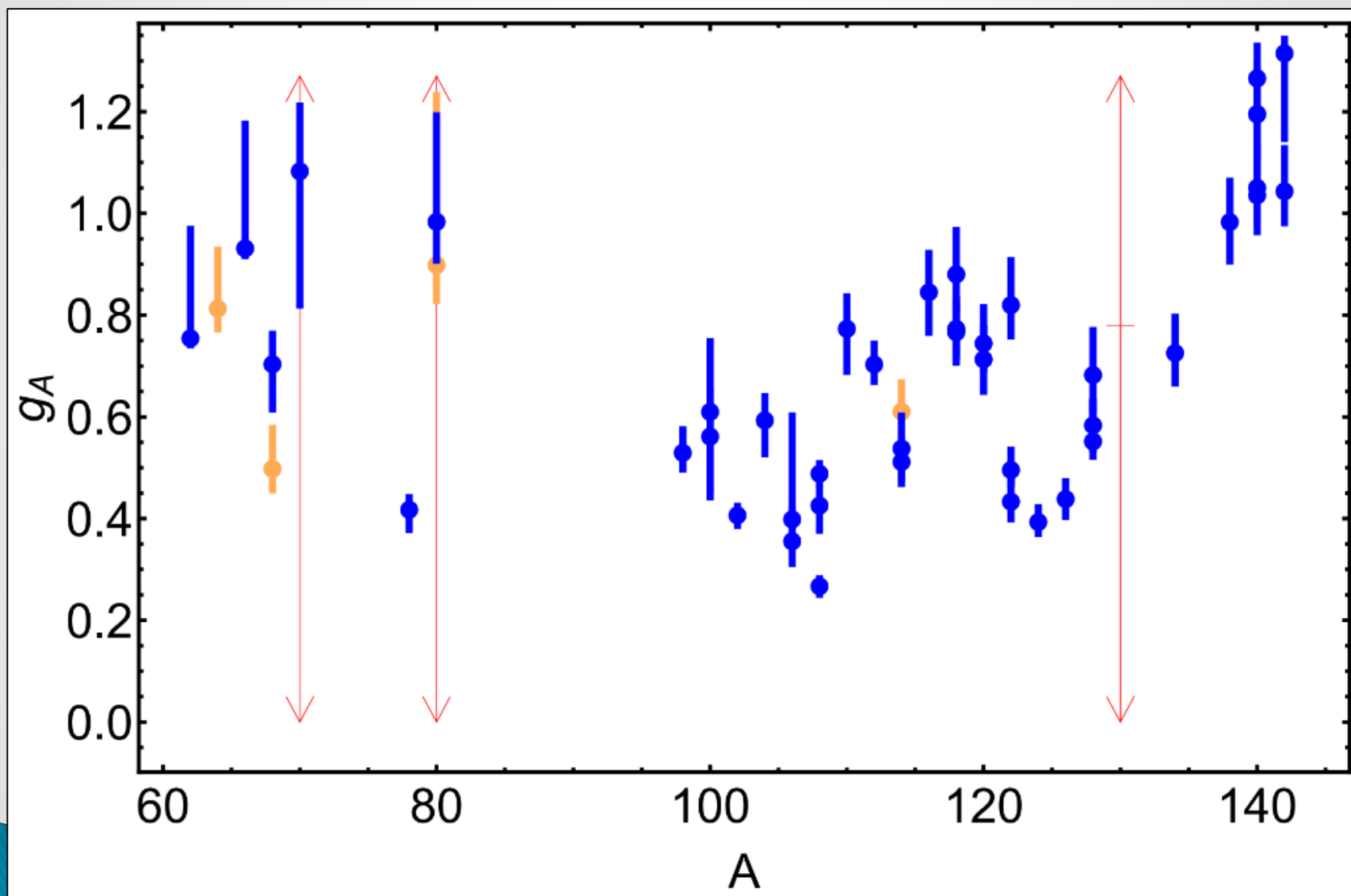


# Isobaric Multiplets

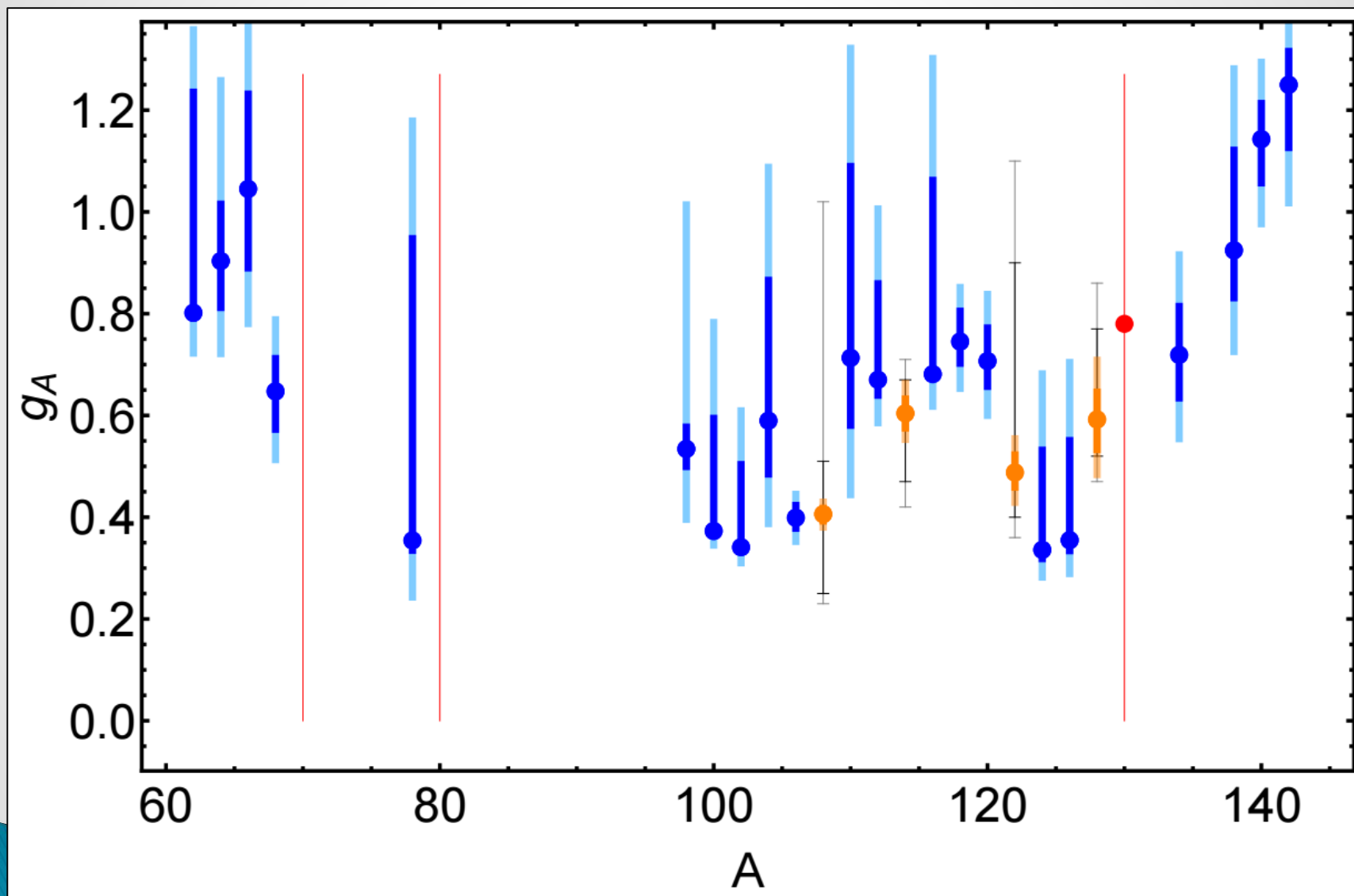
A	Z <sub>0</sub>	Multiplet	dof	$g_A^{\text{fit}}$
62	28	28 ← 29 ← 30	3-2	$0.80^{+0.43}_{-0.01}$
64	28	28 ← 29 → 30	3-2	$0.90^{+0.11}_{-0.09}$
66	28	28 → 29 → 30	3-2	$1.00^{+0.19}_{-0.16}$
68	29	29 → 30 ← 31 ← 32	3-3	$0.65^{+0.06}_{-0.07}$
70	29	29 → <u>30</u> ← 31 → 32	3-3	-
78	34	34 ← 35 → 36	3-2	$0.35^{+0.59}_{-0.02}$
80	33	33 → <u>34</u> ← 35 → 36 ← 37	3-4	<b>1.40</b>
98	39	39 → 40 → 41	2-2	$0.53^{+0.04}_{-0.03}$
100	41	41 → <u>42</u> ← 43 → 44	3-3	$0.37^{+0.22}_{-0.00}$
102	42	42 → 43 → 44	3-2	$0.34^{+0.16}_{-0.00}$
104	44	<u>44</u> ← 45 → 46	3-2	$0.59^{+0.28}_{-0.10}$
106	45	45 → 46 ← 47 → 48	3-3	$0.40^{+0.02}_{-0.02}$
108	44	44 → 45 → 46 ← 47 → 48	4-4	$0.41^{+0.01}_{-0.01}$
110	46	<u>46</u> ← 47 → 48	3-2	$0.71^{+0.38}_{-0.13}$
112	48	48 ← 49 → 50	3-2	$0.67^{+0.19}_{-0.03}$
114	46	46 → 47 → <u>48</u> ← 49 → 50	4-4	$0.60^{+0.03}_{-0.03}$
116	48	<u>48</u> ← 49 → 50	3-2	$0.68^{+0.38}_{-0.01}$
118	48	48 → 49 → 50 ← 51 ← 52	4-4	$0.75^{+0.06}_{-0.04}$
120	48	48 → 49 → 50 ← 51	3-3	$0.71^{+0.06}_{-0.05}$
122	48	48 → 49 → <u>50</u>   52 ← 53 ← 54 ← 55	5-5	$0.49^{+0.03}_{-0.03}$
124	54	54 ← 55 ← 56	3-2	$0.34^{+0.20}_{-0.02}$
126	54	54 ← 55 ← 56	3-2	$0.35^{+0.20}_{-0.02}$
128	52	<u>52</u> ← 53 → 54 ← 55 ← 56	4-4	$0.59^{+0.05}_{-0.06}$
130	54	54 ← 55 → 56	3-2	<b>0.78</b>
134	56	56 ← 57 ← 58	3-2	$0.72^{+0.10}_{-0.08}$
138	58	58 ← 59 ← 60	3-2	$0.92^{+0.20}_{-0.09}$
140	58	58 ← 59 ← 60 ← 61 ← 62 ← 63 ← 64	5-6	$1.10^{+0.07}_{-0.09}$
142	60	60 ← 61 ← 62 ← 63	3-3	$1.20^{+0.07}_{-0.12}$



# Full Results - Triplets



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# Quenching of $g_A$ ?

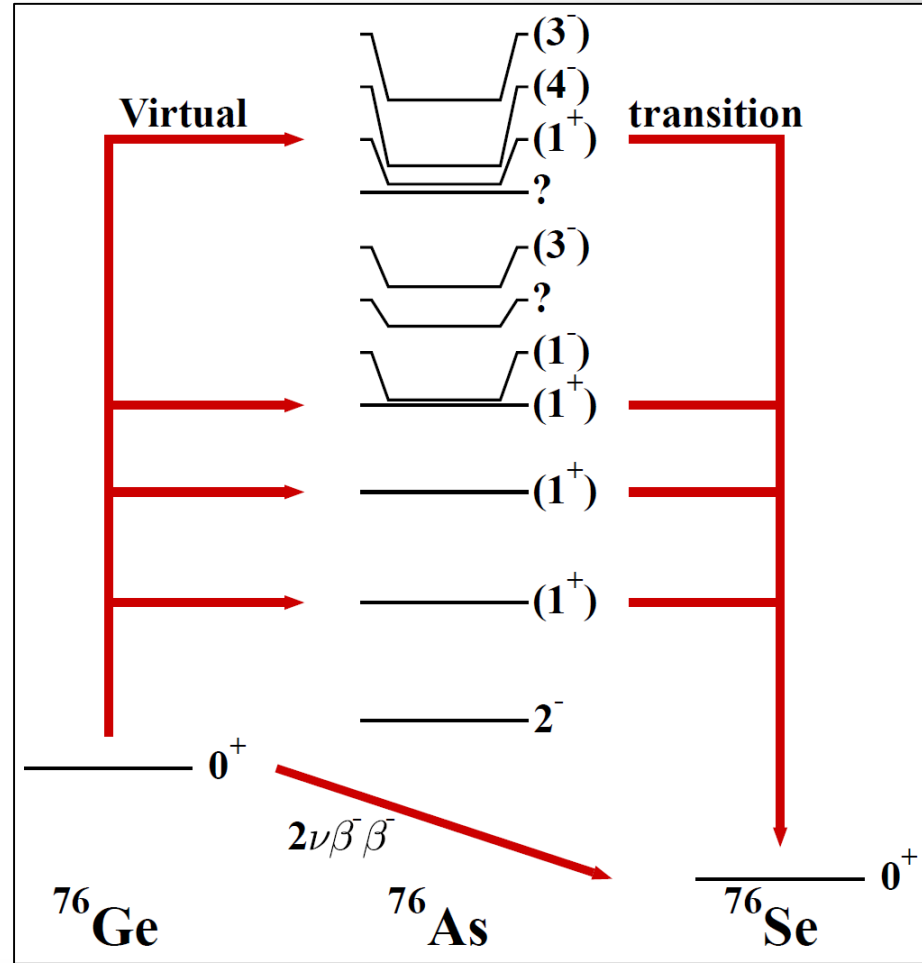
- ▶ Single beta / EC /  $2\nu\beta\beta$   
analysis relevant for  $0\nu\beta\beta$ ?

# Quenching of $g_A$ ?

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Unclear!

- ▶ Processes different at nucleon level
- ▶ Probing different transitions
- ▶ Incorporate more experimental information
  - Higher, forbidden beta decays
  - Charge exchange reactions
  - Muon capture



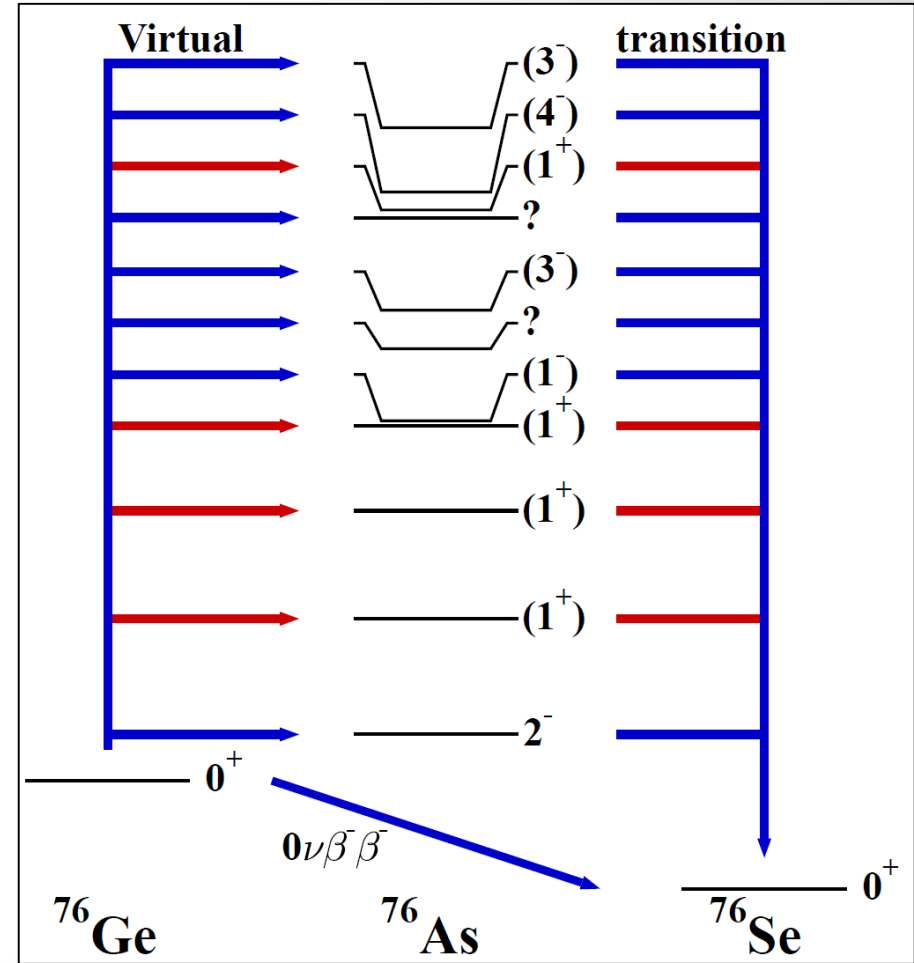
Suhonen '16

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Suhonen '16

- ▶ **Interpretation of  $0\nu\beta\beta$  requires precise NMEs**
  - Estimate of experimental sensitivity
  - Determination of  $0\nu\beta\beta$  mass or falsification of Majorana scenario
- ▶ **Need to determine magnitude of  $g_A$  quenching in  $0\nu\beta\beta$** 
  - At least: additional source of uncertainty
  - Potentially: reduced sensitivity
- ▶ **Solution likely requires concerted effort**
  - Theoretical improvements, e.g. 2-body currents at higher momentum transfer [Menendez, Gazit, Schwenk, Phys. Rev. Lett. 107 (2011) 062501]
  - Experimental probes, e.g. NUMEN
  - Unbiased confrontation of theory with experiment
- ▶ **Given analysis example of consistent fit**
  - Full theory parameter variation against combined experimental data