Single Particle Levels and ββ decay matrix elements in the Interacting Boson Model

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The Interacting Boson Model

$$H_{IBM2} = H_{\pi} + H_{\nu} + \kappa Q_{\pi}^{(2)} \cdot Q_{\nu}^{(2)} + M_{\pi\nu} \left(\xi_{1}, \xi_{2}, \xi_{3}\right)$$

$$H_{\rho} = \epsilon \hat{n}_{d_{\rho}} + \sum_{L=0,2,4} c_{L}^{\rho} \left(d_{\rho}^{\dagger} d_{\rho}^{\dagger}\right)^{(L)} \cdot \left(\tilde{d}_{\rho} \tilde{d}_{\rho}\right)^{(L)}$$

$$Q_{\rho}^{(2)} = \left(s_{\rho}^{\dagger} \times \tilde{d}_{\rho} + d_{\rho}^{\dagger} \times \tilde{s}_{\rho}\right)^{(2)} + \chi_{\rho} \left(d_{\rho}^{\dagger} \times \tilde{d}_{\rho}\right)^{(2)}$$



IBM-2 and Shell Model



Structure of the S and D pairs

BCS method + NOA approximation

$$\frac{\alpha_{j}}{\sqrt{\Omega_{e}}} = \frac{v_{j}}{\sqrt{N_{B}}}, \quad \Omega_{e} = \sum_{j} \alpha_{j}^{2} \left(j + \frac{1}{2} \right)$$
$$\beta_{jj'} = \frac{\langle j \| r^{2} Y^{(2)} \| j' \rangle (u_{j} v_{j'} + v_{j} u_{j'})}{K_{\beta}}, \quad \sum_{j \le j'} \beta_{jj'}^{2} = 1$$

SDI method

$$\begin{vmatrix} 0_1^+ \rangle = \sum_j A_j \ |j^2; 0\rangle \to \alpha_j = \sqrt{\frac{\sum_j (j+\frac{1}{2})}{j+\frac{1}{2}}} A_j \\ \begin{vmatrix} 2_1^+ \rangle = \sum_{j \le j'} B_{jj'} \ |jj'; 2\rangle \to \beta_{jj'} = B_{jj'} \end{vmatrix}$$

Occupancies in the IBM-2

PHYSICAL REVIEW C 94, 034320 (2016)

Occupation probabilities of single particle levels using the microscopic interacting boson model: Application to some nuclei of interest in neutrinoless double- β decay

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$$\hat{n}_{j} = C_{j}^{\dagger} \cdot \tilde{C}_{j} \mapsto \hat{n}_{j}^{B} = A_{j}s^{\dagger} \cdot \tilde{s} + B_{j}d^{\dagger} \cdot \tilde{d}$$

$$A_{j} = \langle 2N, 0, 0 || \hat{n}_{j} || 2N, 0, 0 \rangle$$

$$B_{j} = \frac{1}{\sqrt{5}} \langle 2N, 2, 2 || \hat{n}_{j} || 2N, 2, 2 \rangle$$

$$- (1 - \frac{1}{N}) \langle 2N, 0, 0 || \hat{n}_{j} || 2N, 0, 0 \rangle$$

Updated Single Particle Energies

Occupancies for A = 76



Occupancies for A = 100



Occupancies for A = 130



$$h_X^{F,GT,T} = \sum_{\substack{j_\pi j'_\pi \\ j_\nu j'_\nu \\ J}} -\frac{1}{4} (-1)^J G_X^{F,GT,T} (j_\pi j'_\pi j_\nu j'_\nu; J) \\ \times \sqrt{1 + (-1)^J \delta_{j_\pi j'_\pi}} \sqrt{1 + (-1)^J \delta_{j_\nu j'_\nu}} \\ \times \sqrt{1 + (-1)^J \delta_{j_\pi j'_\pi}} \sqrt{1 + (-1)^J \delta_{j_\nu j'_\nu}} \\ \times \left(C^{\dagger}_{j_\pi} \times C^{\dagger}_{j'_\pi} \right)^{(J)} \cdot \left(\tilde{C}_{j_\nu} \times \tilde{C}_{j'_\nu} \right)^{(J)}$$

$$\begin{pmatrix} C_{j_{\pi}}^{\dagger} \times C_{j_{\pi}}^{\dagger} \end{pmatrix}^{(0)} \longmapsto A_{j_{\pi}} s_{\pi}^{\dagger} + A_{j_{\pi}}' s_{\pi}^{\dagger} \left(d_{\pi}^{\dagger} \tilde{d}_{\pi} \right)^{(0)} + \dots \\ \left(C_{j_{\pi}}^{\dagger} \times C_{j_{\pi}}^{\dagger} \right)^{(2)} \longmapsto B_{j_{\pi}j_{\pi}'} d_{\pi}^{\dagger} + C_{j_{\pi}j_{\pi}'} s_{\pi}^{\dagger} \left(s_{\pi}^{\dagger} \tilde{d}_{\pi} \right)^{(2)} + \dots \\ \left(\tilde{C}_{j_{\nu}} \times \tilde{C}_{j_{\nu}} \right)^{(0)} \longmapsto \tilde{A}_{j_{\nu}} \tilde{s}_{\nu} + \tilde{A}_{j_{\nu}}' \tilde{s}_{\nu} \left(d_{\nu}^{\dagger} \tilde{d}_{\nu} \right)^{(0)} + \dots \\ \left(\tilde{C}_{j_{\nu}} \times \tilde{C}_{j_{\nu}'} \right)^{(2)} \longmapsto \tilde{B}_{j_{\nu}j_{\nu}'} \tilde{d}_{\nu} + \tilde{C}_{j_{\nu}j_{\nu}'} \left(d_{\nu}^{\dagger} \tilde{s}_{\nu} \right)^{(2)} \tilde{s}_{\nu} + \dots$$

$$h_{X}^{F,GT,T} = \sum_{\substack{j_{\pi}j'_{\pi} \\ j_{\nu}j'_{\nu} \\ J}} -\frac{1}{4} (-1)^{J} G_{X}^{F,GT,T} (j_{\pi}j'_{\pi}j_{\nu}j'_{\nu};J) \\ \times \sqrt{1 + (-1)^{J} \delta_{j_{\pi}j'_{\pi}}} \sqrt{1 + (-1)^{J} \delta_{j_{\nu}j'_{\nu}}} \\ \times \left(C^{\dagger}_{j_{\pi}} \times C^{\dagger}_{j'_{\pi}} \right)^{(J)} \cdot \left(\tilde{C}_{j_{\nu}} \times \tilde{C}_{j'_{\nu}} \right)^{(J)}$$





$$\begin{pmatrix} C_{j_{\pi}}^{\dagger} \times C_{j_{\pi}}^{\dagger} \end{pmatrix} \longmapsto B_{j_{\pi}j_{\pi}'} d_{\pi}^{\dagger} + C_{j_{\pi}j_{\pi}'} s_{\pi}^{\dagger} \left(s_{\pi}^{\dagger} d_{\pi} \right) + \dots$$

$$\begin{pmatrix} \tilde{C}_{j_{\nu}} \times \tilde{C}_{j_{\nu}} \end{pmatrix}^{(0)} \longmapsto \tilde{A}_{j_{\nu}} \tilde{s}_{\nu} + \tilde{A}_{j_{\nu}}' \tilde{s}_{\nu} \left(d_{\nu}^{\dagger} \tilde{d}_{\nu} \right)^{(0)} + \dots$$

$$\begin{pmatrix} \tilde{C}_{j_{\nu}} \times \tilde{C}_{j_{\nu}'} \end{pmatrix}^{(2)} \longmapsto \tilde{B}_{j_{\nu}j_{\nu}'} \tilde{d}_{\nu} + \tilde{C}_{j_{\nu}j_{\nu}'} \left(d_{\nu}^{\dagger} \tilde{s}_{\nu} \right)^{(2)} \tilde{s}_{\nu} + \dots$$

$$h_X^{F,GT,T} = \sum_{\substack{j_\pi j'_\pi \\ j_\nu j'_\nu \\ J}} -\frac{1}{4} (-1)^J G_X^{F,GT,T} (j_\pi j'_\pi j_\nu j'_\nu; J) \\ \times \sqrt{1 + (-1)^J \delta_{j_\pi j'_\pi}} \sqrt{1 + (-1)^J \delta_{j_\nu j'_\nu}} \\ \times \sqrt{1 + (-1)^J \delta_{j_\pi j'_\pi}} \sqrt{1 + (-1)^J \delta_{j_\nu j'_\nu}} \\ \times \left(C^{\dagger}_{j_\pi} \times C^{\dagger}_{j'_\pi} \right)^{(J)} \cdot \left(\tilde{C}_{j_\nu} \times \tilde{C}_{j'_\nu} \right)^{(J)}$$

$$\begin{pmatrix}
\begin{pmatrix}
C_{j_{\pi}}^{\dagger} \times C_{j_{\pi}}^{\dagger}
\end{pmatrix}^{(0)} & \longmapsto & A_{j_{\pi}}s_{\pi}^{\dagger} + A'_{j_{\pi}}s_{\pi}^{\dagger}\left(d_{\pi}^{\dagger}\tilde{d}_{\pi}\right)^{(0)} + \dots \\
\begin{pmatrix}
C_{j_{\pi}}^{\dagger} \times C_{j'_{\pi}}^{\dagger}
\end{pmatrix}^{(2)} & \longmapsto & B_{j_{\pi}j'_{\pi}}d_{\pi}^{\dagger} + C_{j_{\pi}j'_{\pi}}s_{\pi}^{\dagger}\left(s_{\pi}^{\dagger}\tilde{d}_{\pi}\right)^{(2)} + \dots \\
\begin{pmatrix}
\tilde{C}_{j_{\nu}} \times \tilde{C}_{j_{\nu}}
\end{pmatrix}^{(0)} & \longmapsto & \tilde{A}_{j_{\nu}}\tilde{s}_{\nu} + \tilde{A}'_{j_{\nu}}\tilde{s}_{\nu}\left(d_{\nu}^{\dagger}\tilde{d}_{\nu}\right)^{(0)} + \dots \\
\begin{pmatrix}
\tilde{C}_{j_{\nu}} \times \tilde{C}_{j'_{\nu}}
\end{pmatrix}^{(2)} & \longmapsto & \tilde{B}_{j_{\nu}j'_{\nu}}\tilde{d}_{\nu} + \tilde{C}_{j_{\nu}j'_{\nu}}\left(d_{\nu}^{\dagger}\tilde{s}_{\nu}\right)^{(2)}\tilde{s}_{\nu} + \dots
\end{cases}$$

IBM-2 DBD calculation



Neutrinoless DBD, light v



Neutrinoless DBD, heavy v



Two-neutrino DBD



Changes in the S.P.E.



Occupations for ¹³⁰Te



Evolution of the NME for A = 130

A = 130



Occupations for 150Nd

Evolution of the NME for A = 150

A = 150

Summary

- Occupancies calculated with updated single particle energies reproduce data reasonably.
- DBD NMEs increase with the updated single particle energies for some cases.
- The increase is enhanced when nuclei move away from closed shells.

Thanks

Effects of the mapping coefficients

$$B_{jj'} = -\sqrt{5(1+\delta_{jj'})} \frac{\eta_{n+2,2,2}^2(j'j)}{\eta_{n,0,0}\eta_{n+2,2,2}} \beta_{j'j}$$

$$\eta_{n,0,0}^2 = \left(\frac{n}{2}!\right)^2 \sum_{m_1...m_k} \left\{ \prod_{i=1}^k \alpha_{j_i}^{2m_i} \left(\begin{array}{c} \Omega_{j_i} \\ m_i \end{array}\right) \right\}$$

$$\eta_{n,2,2}^2 = \sum_{j \le j'} \beta_{jj'}^2 \eta_{n,2,2}^2(jj')$$

$$\eta_{n,2,2}^2(jj') = \sum_{\substack{p=0\\p=0}}^{\frac{n}{2}-1} \left[\frac{(\frac{n}{2}-1)!}{p!}\right]^2 (-1)^{\frac{n}{2}-1-p} \eta_{2p,0,0}^2$$

$$\times \sum_{\substack{q=0\\q=0}}^{\frac{n}{2}-1-p} \alpha_j^{n-2-2p-2q} \alpha_{j'}^{2q}$$